

AN ANALYTICAL STUDY OF SPACE-TIME FRACTIONAL ORDER GAS DYNAMIC EQUATIONS

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Abstract

In this article, the Sumudu transform with iterative method is implemented to obtain approximate analytical solutions in series form to non-linear homogeneous and non-homogeneous space-time fractional gas dynamic equations. The fractional derivatives presented here are in the Caputo sense. Furthermore, the findings of this study are graphically represented and the solution graphs demonstrate a strong connection between the approximate and exact solutions.

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1 Introduction

Fractional calculus is a branch of applied mathematics that is extremely useful in a variety of fields of research [13, 21]. The fractional differential equations have sparked the interest of a vast scope of researchers working on a variety of applications [2, 3, 10, 28, 31]. Many efforts have been made to develop analytical and numerical approaches for solving differential equations of fractional order, such as the homotopy analysis method (*HAM*) [17], the q -homotopy analysis method (q -*HAM*) [14], the optimal q -homotopy analysis method (Oq -*HAM*) [32], the homotopy analysis transform method (*HATM*) [30], the adomian decomposition method (*ADM*) [12], the Laplace decomposition method (*LDM*) [18], the homotopy perturbation method (*HPM*) [19], the homotopy perturbation transform method (*HPTM*) [20, 23], and so on.

In 2006, Daftardar-Gejji and Jafari [8, 15] proposed an iterative method for numerically solving nonlinear functional equations. Since then, the iterative technique has been used to solve a wide variety of nonlinear differential equations of integer and fractional order [5] as well as fractional boundary value problems [7]. Recently, Wang and Liu [33] introduced the Sumudu transform iterative method (*STIM*) by combining the Sumudu transform with an iterative technique to determine approximate analytical solutions of time-fractional Cauchy reaction diffusion equations. The Sumudu transform iterative technique has been used successfully to solve a variety of time and space fractional partial differential equations and related systems [22], as well as the random component time-fractional Klein-Gordon equation [27].

In this work, we consider the non-linear homogeneous and non-homogeneous fractional gas dynamic equations with space and time fractional derivatives as follows

(i) The non-linear homogeneous space-time-fractional gas dynamic partial differential equation of the form

$$(1.1) \quad D_t^\alpha u(x, t) + \frac{1}{2} D_x^\beta u^2(x, t) - u(x, t)(1 - u(x, t)) = 0, \quad 0 < \alpha, \beta \leq 1,$$

$$(1.2) \quad u(x, 0) = g(x),$$

(ii) The non-linear non-homogeneous space-time-fractional gas dynamic partial differential equation of the form

$$(1.3) \quad D_t^\alpha u(x, t) + \frac{1}{2} D_x^\beta u^2(x, t) - u(x, t)(1 - u(x, t)) = f(x, t), \quad 0 < \alpha, \beta \leq 1,$$

$$(1.4) \quad u(x, 0) = g(x),$$

where α and β are the parameters that describe the order of the time-fractional and space-fractional derivatives, respectively. Also, $u(x, t)$ is the probability density function and f is a known analytic function.

2 Preliminaries and Basic Definitions

This section introduces some fundamental definitions, notations, and properties of fractional calculus utilizing Sumudu transform theory, which will be applied later in this paper.

Definition 2.1. The Caputo fractional derivative of a function $u(x, t)$ is defined as [28, 31]

$$(2.1) \quad D_t^\alpha u(x, t) = \frac{1}{\Gamma(m-\alpha)} \int_0^t (t-\eta)^{m-\alpha-1} u^{(m)}(x, \eta) d\eta, \quad m-1 < \alpha \leq m, m \in \mathbb{N}.$$

Definition 2.2. The Sumudu transform is defined over the set of functions

$$\left\{ f(t) \mid \exists M, \rho_1 > 0, \rho_2 > 0, |f(t)| < M e^{t/\rho_j} \text{ if } t \in (-1)^j \times [0, \infty), j = 1, 2 \right\}$$

by the following formula [4, 34]

$$(2.2) \quad S[f(t)] = F(\omega) = \int_0^\infty e^{-t} f(\omega t) dt, \quad \omega \in (-\rho_1, \rho_2).$$

Definition 2.3. The Sumudu transform of Caputo fractional derivative is defined in the following manner [9, 33]

$$(2.3) \quad S[D_t^\alpha u(x, t)] = \omega^{-\alpha} S[u(x, t)] - \sum_{k=0}^{m-1} \omega^{-\alpha+k} u^{(k)}(x, 0), \quad m-1 < \alpha \leq m, m \in \mathbb{N},$$

where $u^{(k)}(x, 0)$ is the k -order derivative of $u(x, t)$ with respect to t at $t = 0$.

Definition 2.4. The Mittag-Leffler function, a generalization of the exponential function, is defined as follows [28, 31]

$$(2.4) \quad E_\alpha(z) = \sum_{n=0}^{\infty} \frac{z^n}{\Gamma(\alpha n + 1)}, \quad \alpha \in \mathbb{C}, \operatorname{Re}(\alpha) > 0.$$

A further generalization of equation (2.4) is as follows [35]

$$(2.5) \quad E_{\alpha, \beta}(z) = \sum_{n=0}^{\infty} \frac{z^n}{\Gamma(\alpha n + \beta)}, \quad \alpha, \beta \in \mathbb{C}, \operatorname{Re}(\alpha) > 0, \operatorname{Re}(\beta) > 0,$$

where $\Gamma(\cdot)$ is the well-known Gamma function.

3 Basic Idea of Sumudu Transform Iterative Method

To explain the basic idea of the Sumudu transform iterative technique [33], we take the following space and time general fractional partial differential equation having the prescribed initial conditions may be written in the form of an operator as

$$(3.1) \quad D_t^\alpha u(x, t) = F[x, u(x, t), D_x^\beta u(x, t), \dots, D_x^{l\beta} u(x, t)],$$

$$l-1 < \alpha \leq l, \quad m-1 < \beta \leq m; \quad l, m \in \mathbb{N}$$

$$(3.2) \quad u^{(k)}(x, 0) = h_k(x), \quad k = 0, 1, 2, \dots, n-1,$$

where $D_t^\alpha u(x, t)$ and $D_x^\beta u(x, t)$ are the Caputo fractional derivatives of order α , $l-1 < \alpha \leq l$ and β , $m-1 < \beta \leq m$, respectively, defined by the equation (2.1), $F[x, u, D_x^\beta u, \dots, D_x^{l\beta} u]$ is a linear/non-linear operator and $u = u(x, t)$ is the unknown function and fractional derivative $D_x^{l\beta} u(x, t)$, $l \in \mathbb{N}$ is taken as the sequential fractional derivative [28] that is

$$(3.3) \quad D_x^{l\beta} u = D_x^\beta D_x^\beta, \dots, D_x^\beta u \quad (l \text{ times}).$$

Applying the Sumudu transform on both sides of equation (3.1), we have

$$(3.4) \quad S[D_t^\alpha u(x, t)] = S[F(x, u(x, t), D_x^\beta u(x, t), \dots, D_x^{l\beta} u(x, t))].$$

Using the differentiation property of the Sumudu transform, we get

$$(3.5) \quad S[u(x, t)] = \omega^\alpha \sum_{k=0}^{m-1} [\omega^{-\alpha+k} u^{(k)}(x, 0)] + \omega^\alpha S[F(x, u, D_x^\beta u, \dots, D_x^{l\beta} u)].$$

On taking inverse Sumudu transform of equation (3.5), we have

$$(3.6) \quad u(x, t) = S^{-1} \left[\omega^\alpha \sum_{k=0}^{m-1} [\omega^{-\alpha+k} u^{(k)}(x, 0)] \right] + S^{-1} \left[\omega^\alpha S[F(x, u, D_x^\beta u, \dots, D_x^{l\beta} u)] \right].$$

Equation (3.6) may be written as

$$(3.7) \quad u(x, t) = f(x, t) + N\left(x, u, D_x^\beta u, \dots, D_x^{l\beta} u\right),$$

where

$$(3.8) \quad f(x, t) = S^{-1} \left[\omega^\alpha \sum_{k=0}^{m-1} \left[\omega^{-\alpha+k} u^{(k)}(x, 0) \right] \right],$$

$$(3.9) \quad N\left(x, u, D_x^\beta u, \dots, D_x^{l\beta} u\right) = S^{-1} \left[\omega^\alpha S \left[F\left(x, u, D_x^\beta u, \dots, D_x^{l\beta} u\right) \right] \right].$$

Here N is a linear/nonlinear operator and f is a known function .

Furthermore, we employ the iterative method proposed by Daftardar-Gejji and Jafari [8], which represents a solution in an infinite series of components as

$$(3.10) \quad u(x, t) = \sum_{i=0}^{\infty} u_i(x, t).$$

The operator N is decomposed as follows

$$(3.11) \quad N\left(x, \sum_{i=0}^{\infty} u_i, D_x^\beta \left(\sum_{i=0}^{\infty} u_i\right), \dots, D_x^{l\beta} \left(\sum_{i=0}^{\infty} u_i\right)\right) = N\left(x, u_0, D_x^\beta u_0, \dots, D_x^{l\beta} u_0\right) \\ + \sum_{j=1}^{\infty} \left[N\left(x, \sum_{i=0}^j u_i, D_x^\beta \left(\sum_{i=0}^j u_i\right), \dots, D_x^{l\beta} \left(\sum_{i=0}^j u_i\right)\right) \right] \\ - \sum_{j=1}^{\infty} \left[N\left(x, \sum_{i=0}^{j-1} u_i, D_x^\beta \left(\sum_{i=0}^{j-1} u_i\right), \dots, D_x^{l\beta} \left(\sum_{i=0}^{j-1} u_i\right)\right) \right],$$

$$(3.12) \quad S^{-1} \left[\omega^\alpha S \left[F\left(x, \sum_{i=0}^{\infty} u_i, D_x^\beta \left(\sum_{i=0}^{\infty} u_i\right), \dots, D_x^{l\beta} \left(\sum_{i=0}^{\infty} u_i\right)\right) \right] \right] \\ = S^{-1} \left[\omega^\alpha S \left[F\left(x, u_0, D_x^\beta u_0, \dots, D_x^{l\beta} u_0\right) \right] \right] \\ + \sum_{j=0}^{\infty} \left[S^{-1} \left[\omega^\alpha S \left[F\left(x, \sum_{i=0}^j u_i, D_x^\beta \left(\sum_{i=0}^j u_i\right), \dots, D_x^{l\beta} \left(\sum_{i=0}^j u_i\right)\right) \right] \right] \right] \\ - \sum_{j=0}^{\infty} \left[S^{-1} \left[\omega^\alpha S \left[F\left(x, \sum_{i=0}^{j-1} u_i, D_x^\beta \left(\sum_{i=0}^{j-1} u_i\right), \dots, D_x^{l\beta} \left(\sum_{i=0}^{j-1} u_i\right)\right) \right] \right] \right].$$

Using equations (3.10) to (3.12) in equation (3.7), we obtain

$$(3.13) \quad \sum_{i=0}^{\infty} u_i(x, t) = S^{-1} \left[\omega^\alpha \sum_{k=0}^{m-1} \left(\omega^{-\alpha+k} u^{(k)}(x, 0) \right) \right] \\ + S^{-1} \left[\omega^\alpha S \left[F\left(x, u_0, D_x^\beta u_0, \dots, D_x^{l\beta} u_0\right) \right] \right] \\ + \sum_{j=0}^{\infty} \left[S^{-1} \left[\omega^\alpha S \left[F\left(x, \sum_{i=0}^j u_i, D_x^\beta \left(\sum_{i=0}^j u_i\right), \dots, D_x^{l\beta} \left(\sum_{i=0}^j u_i\right)\right) \right] \right] \right] \\ - \sum_{j=0}^{\infty} \left[S^{-1} \left[\omega^\alpha S \left[F\left(x, \sum_{i=0}^{j-1} u_i, D_x^\beta \left(\sum_{i=0}^{j-1} u_i\right), \dots, D_x^{l\beta} \left(\sum_{i=0}^{j-1} u_i\right)\right) \right] \right] \right].$$

The recurrence relations have been defined as follows

$$(3.14) \quad u_0(x, t) = S^{-1} \left[\omega^\alpha \sum_{k=0}^{m-1} \left(\omega^{-\alpha+k} u^{(k)}(x, 0) \right) \right],$$

$$(3.15) \quad u_1(x, t) = S^{-1} \left[\omega^\alpha S \left[F\left(x, u_0, D_x^\beta u_0, \dots, D_x^{l\beta} u_0\right) \right] \right],$$

$$(3.16) \quad u_{r+1}(x, t) = S^{-1} \left[\omega^\alpha S \left[F \left(x, \sum_{i=0}^r u_i, D_x^\beta \left(\sum_{i=0}^r u_i \right), \dots, D_x^{l\beta} \left(\sum_{i=0}^r u_i \right) \right) \right] \right. \\ \left. - S^{-1} \left[\omega^\alpha S \left[F \left(x, \sum_{i=0}^{r-1} u_i, D_x^\beta \left(\sum_{i=0}^{r-1} u_i \right), \dots, D_x^{l\beta} \left(\sum_{i=0}^{r-1} u_i \right) \right) \right] \right], r \geq 1.$$

Therefore, the approximate analytical solution of equations (3.1) and (3.2) in truncated series form is given by

$$(3.17) \quad u(x, t) \cong \lim_{N \rightarrow \infty} \sum_{m=0}^N u_m(x, t).$$

In general, the solutions in the above series converge quickly. The classical approach to the convergence of this type of series has been presented by Bhalekar and Daftardar-Gejji [6] and Daftardar-Gejji and Jafari [8].

4 Solution of the Space-Time Fractional Gas Dynamic Equations

In this section, we make an attempt to solve non-linear homogeneous and non-homogeneous space-time fractional gas dynamic equations by means of the Sumudu transform iterative method.

Example 4.1. Consider the following non-linear homogeneous space-time fractional gas dynamic equation [30, 32]

$$(4.1) \quad D_t^\alpha u(x, t) + \frac{1}{2} D_x^\beta u^2(x, t) - u(x, t)(1 - u(x, t)) = 0, \quad t > 0, \quad 0 < \alpha, \beta \leq 1,$$

with the initial condition

$$(4.2) \quad u(x, 0) = e^{-x}.$$

Taking the Sumudu transform on the both sides of equation (4.1), and making use of the result given by equation (4.2), we have

$$(4.3) \quad S[u(x, t)] = e^{-x} + \omega^\alpha S \left[-\frac{1}{2} \frac{\partial^\beta u^2(x, t)}{\partial x^\beta} + u(x, t)(1 - u(x, t)) \right].$$

On taking inverse Sumudu transform of equation (4.3), we get

$$(4.4) \quad u(x, t) = e^{-x} + S^{-1} \left[\omega^\alpha S \left[-\frac{1}{2} \frac{\partial^\beta u^2(x, t)}{\partial x^\beta} + u(x, t)(1 - u(x, t)) \right] \right].$$

Substituting the results from equations (3.10) to (3.12) in the equation (4.4) and applying the equations (3.14) to (3.16), we determine the components of the solution as follows

$$(4.5) \quad u_0(x, t) = u(x, 0) = e^{-x},$$

$$(4.6) \quad u_1(x, t) = S^{-1} \left[\omega^\alpha S \left[-\frac{1}{2} \frac{\partial^\beta u_0^2}{\partial x^\beta} + u_0(1 - u_0) \right] \right] \\ = -e^{-x} \left[\frac{t^\alpha}{\Gamma(\alpha + 1)} \right] \left(2^{-1+\beta} e^{-x+i\beta\pi} - 1 + e^{-x} \right),$$

$$(4.7) \quad u_2(x, t) = S^{-1} \left[\omega^\alpha S \left[-\frac{1}{2} \frac{\partial^\beta (u_0 + u_1)^2}{\partial x^\beta} + (u_0 + u_1)(1 - (u_0 + u_1)) \right] \right] \\ - S^{-1} \left[\omega^\alpha S \left[-\frac{1}{2} \frac{\partial^\beta u_0^2}{\partial x^\beta} + u_0(1 - u_0) \right] \right] \\ = \left[-e^{-x} \frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} + 2e^{-2x} \frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} \right] \left(2^{-1+\beta} e^{-x+i\beta\pi} - 1 + e^{-x} \right) \\ + \left[\frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} \right] \left(3^\beta e^{-3x+i\beta\pi} - 2^\beta e^{-2x+i\beta\pi} + 2^{-1+\beta} 3^\beta e^{-3x+2i\beta\pi} \right),$$

and so on. The remaining components may be obtained in the same way.

Thus, the approximate analytical solution in the series form can be obtained as

$$(4.8) \quad u(x, t) \cong \lim_{N \rightarrow \infty} \sum_{m=0}^N u_m(x, t) = u_0(x, t) + u_1(x, t) + u_2(x, t) + \dots,$$

$$\begin{aligned}
&= e^{-x} - e^{-x} \left[\frac{t^\alpha}{\Gamma(\alpha + 1)} \right] (2^{-1+\beta} e^{-x+i\beta\pi} - 1 + e^{-x}) \\
&+ \left[-e^{-x} \frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} + 2e^{-2x} \frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} \right] (2^{-1+\beta} e^{-x+i\beta\pi} - 1 + e^{-x}) \\
&+ \left[\frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} \right] (3^\beta e^{-3x+i\beta\pi} - 2^\beta e^{-2x+i\beta\pi} + 2^{-1+\beta} 3^\beta e^{-3x+2i\beta\pi}) + \dots .
\end{aligned}$$

The same result was obtained by Saad *et al.* [32] by using the method of optimal q -HAM.

If we put $\alpha = \beta = 1$, in equation (4.8), we have the result in simple form

$$(4.9) \quad u(x, t) = e^{t-x},$$

which is the exactly the same solution obtained by earlier by Jafari *et al.* [19] by using HPM method.

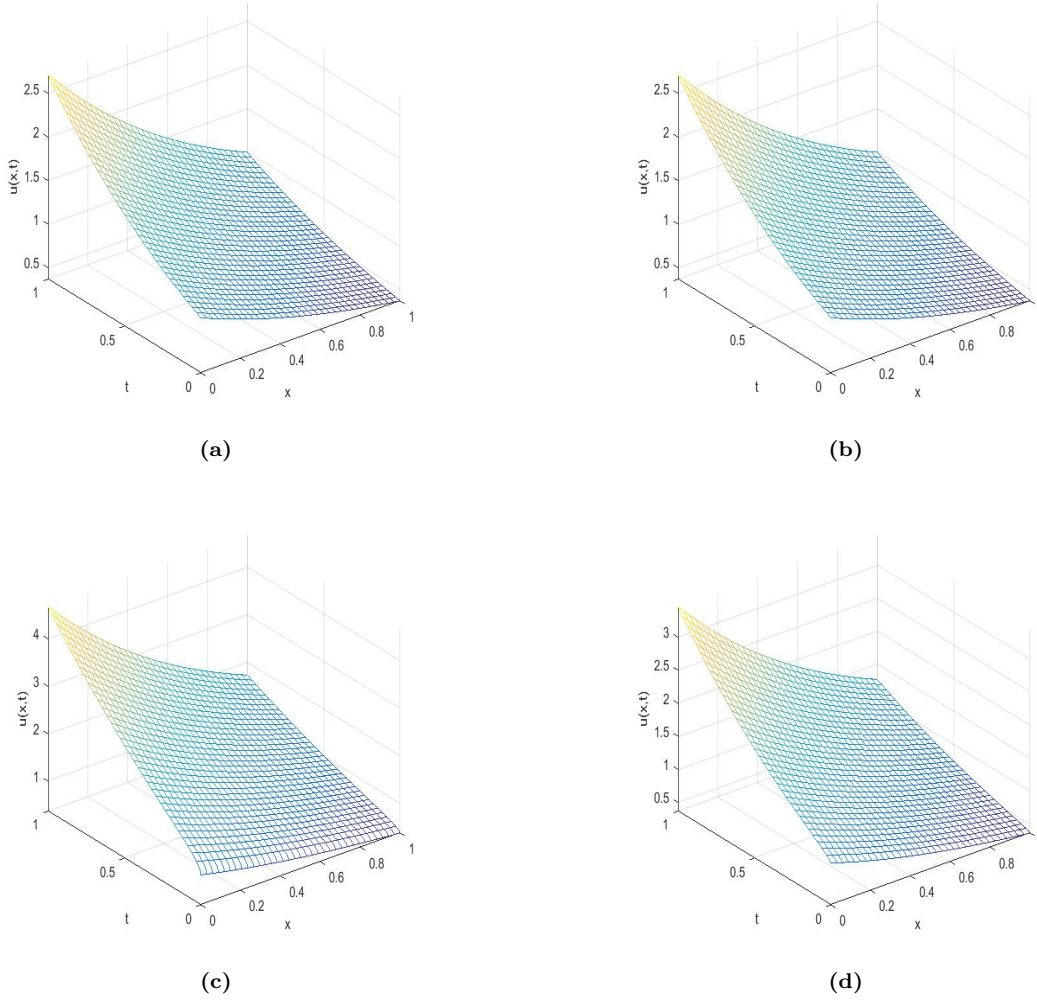


Figure 4.1: Graph of the $u(x, t)$ for Example 4.1, when $\beta = 1$: (a) The exact solution, (b) The approximate solution for $\alpha = 1$, (c) The approximate solution for $\alpha = 0.5$, (d) The approximate solution for $\alpha = 0.75$.

Example 4.2. Consider the following non-linear non-homogeneous space-time fractional gas dynamic equation [32]

$$(4.10) \quad D_t^\alpha u(x, t) + \frac{1}{2} D_x^\beta u^2(x, t) - u(x, t)(1 - u(x, t)) = -e^{t-x}, \quad t > 0, \quad 0 < \alpha, \beta \leq 1,$$

with the initial condition

$$(4.11) \quad u(x, 0) = 1 - e^{-x}.$$

Taking the Sumudu transform on the both sides of equation (4.10), and making use of the result given by equation (4.11), we have

$$(4.12) \quad S[u(x, t)] = 1 - e^{-x} + \omega^\alpha S \left[-\frac{1}{2} \frac{\partial^\beta u^2(x, t)}{\partial x^\beta} + u(x, t)(1 - u(x, t)) - e^{t-x} \right].$$

On taking inverse Sumudu transform of equation (4.12), we get

$$(4.13) \quad u(x, t) = 1 - e^{-x} + S^{-1} \left[\omega^\alpha S \left[-\frac{1}{2} \frac{\partial^\beta u^2(x, t)}{\partial x^\beta} + u(x, t)(1 - u(x, t)) - e^{t-x} \right] \right].$$

Substituting the results from equations (3.10) to (3.12) in the equation (4.13) and applying the equations (3.14) to (3.16), we determine the components of the solution as follows

$$(4.14) \quad u_0(x, t) = u(x, 0) = 1 - e^{-x},$$

$$(4.15) \quad u_1(x, t) = S^{-1} \left[\omega^\alpha S \left[-\frac{1}{2} \frac{\partial^\beta u_0^2}{\partial x^\beta} + u_0(1 - u_0) - e^{t-x} \right] \right] \\ = -e^{-x} \left[\frac{t^\alpha}{\Gamma(\alpha + 1)} \right] \left(-e^{i\beta\pi} + 2^{-1+\beta} e^{-x+i\beta\pi} - 1 + e^{-x} \right) - e^{-x} t^\alpha E_{1, \alpha+1}(t),$$

$$(4.16) \quad u_2(x, t) = S^{-1} \left[\omega^\alpha S \left[-\frac{1}{2} \frac{\partial^\beta (u_0 + u_1)^2}{\partial x^\beta} + (u_0 + u_1)(1 - (u_0 + u_1)) - e^{t-x} \right] \right] \\ - S^{-1} \left[\omega^\alpha S \left[-\frac{1}{2} \frac{\partial^\beta u_0^2}{\partial x^\beta} + u_0(1 - u_0) - e^{t-x} \right] \right] \\ = t^{2\alpha} \left(-e^{i\pi\beta} + 2^{-1+\beta} e^{-x+i\pi\beta} - 1 + e^{-x} \right) \left(-\frac{2e^{-2x}}{\Gamma(2\alpha + 1)} + \frac{e^{-x} t^{2\alpha}}{\Gamma(2\alpha + 1)} \right) \\ - \frac{(3^\beta e^{-3x+i\pi\beta} - 2^\beta e^{-2x+i\pi\beta} + 2^{-1+\beta} 3^\beta e^{-3x+2i\pi\beta} - 2^\beta e^{-2x+2i\pi\beta}) t^{2\alpha}}{\Gamma(2\alpha + 1)} \\ + \frac{(2^\beta e^{-2x+i\pi\beta} - e^{-x+i\pi\beta} + 2^\beta e^{-2x+2i\pi\beta} - e^{-x+2i\pi\beta}) t^{2\alpha}}{\Gamma(2\alpha + 1)} \\ + \left(-2e^{-2x} E_{1, 2\alpha+1}(t) + e^{-x} E_{1, 2\alpha+1}(t) \right. \\ \left. + e^{-x+i\pi\beta} E_{1, 2\alpha+1}(t) - 2^\beta e^{-2x+i\pi\beta} E_{1, 2\alpha+1}(t) \right) t^{2\alpha},$$

and so on. The remaining components may be obtained in the same way.

Thus, the approximate analytical solution in the series form can be obtained as

$$(4.17) \quad u(x, t) \cong \lim_{N \rightarrow \infty} \sum_{m=0}^N u_m(x, t) = u_0(x, t) + u_1(x, t) + u_2(x, t) + \dots, \\ = 1 - e^{-x} - e^{-x} \frac{t^\alpha}{\Gamma(\alpha + 1)} \left(-e^{i\pi\beta} + 2^{-1+\beta} e^{-x+i\pi\beta} - 1 + e^{-x} \right) - e^{-x} t^\alpha E_{1, \alpha+1}(t) \\ + t^{2\alpha} \left(-e^{i\pi\beta} + 2^{-1+\beta} e^{-x+i\pi\beta} - 1 + e^{-x} \right) \left(-\frac{2e^{-2x}}{\Gamma(2\alpha + 1)} + \frac{e^{-x} t^{2\alpha}}{\Gamma(2\alpha + 1)} \right) \\ - \frac{(3^\beta e^{-3x+i\pi\beta} - 2^\beta e^{-2x+i\pi\beta} + 2^{-1+\beta} 3^\beta e^{-3x+2i\pi\beta} - 2^\beta e^{-2x+2i\pi\beta}) t^{2\alpha}}{\Gamma(2\alpha + 1)} \\ + \frac{(2^\beta e^{-2x+i\pi\beta} - e^{-x+i\pi\beta} + 2^\beta e^{-2x+2i\pi\beta} - e^{-x+2i\pi\beta}) t^{2\alpha}}{\Gamma(2\alpha + 1)} \\ + \left(-2e^{-2x} E_{1, 2\alpha+1}(t) + e^{-x} E_{1, 2\alpha+1}(t) \right. \\ \left. + e^{-x+i\pi\beta} E_{1, 2\alpha+1}(t) - 2^\beta e^{-2x+i\pi\beta} E_{1, 2\alpha+1}(t) \right) t^{2\alpha} + \dots .$$

The same result was obtained by Saad *et al.* [32] by using the method of optimal q -HAM. If we put $\alpha = \beta = 1$, in equation (4.17), we have the result in simple form

$$(4.18) \quad u(x, t) = 1 - e^{t-x},$$

which is the exactly the same solution obtained by earlier by Jafari *et al.* [16] by using two-dimensional *DTM* method.

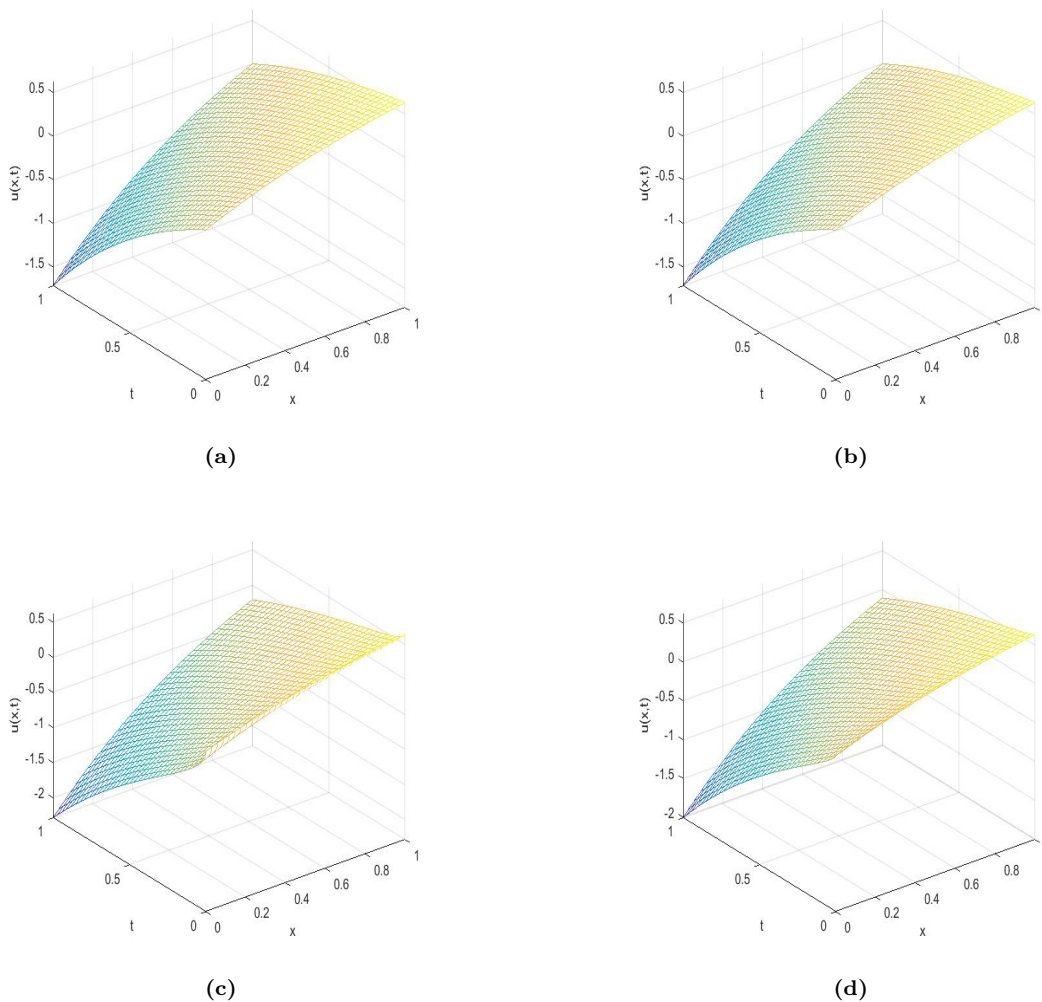


Figure 4.2: Graph of the $u(x, t)$ for Example 4.2, when $\beta = 1$: (a) The exact solution, (b) The approximate solution for $\alpha = 1$, (c) The approximate solution for $\alpha = 0.5$, (d) The approximate solution for $\alpha = 0.75$.

5 Conclusion

In this paper, we have successfully and efficiently applied the Sumudu transform iterative method (*STIM*) to derive the approximate analytical solutions of the non-linear homogeneous and non-homogeneous space-time fractional gas dynamic equations with Caputo fractional derivatives. *STIM* is a hybrid approach of the Sumudu transform and the iterative method. The graphical representation of the obtained solutions was completed successfully by the *MATLAB* software. The analytical results derived from the proposed approach indicate that the method is simple to use and precise.

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